Ion Dependent Computation for Electronic Dielectric Constants of Mixed Crystals in Rocksalt and Zinc blende Structures

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Abstract

Quantum ion dependent theory for electronic dielectric constants of $A^N B^{8-N}$ binary solids has been extended to the II-VI crystals mixed in different proportions. The values of optical dielectric constants (ε_{∞}) are evaluated for a number of mixed crystals, which are found to agree with available experimental values. The results agree well with the predictions from well established Clausius-Mossotti relation also. Thus, it gives a strong support to the validity of ion dependent formulation for mixed crystals.

Key words: Electronic Dielectric Constant (ε_{∞}), Interionic Separation (R), Mixed Crystals

Introduction

he electronic dielectric behaviour of binary

crystals remained a topic of great interest in last few decades as it is a subject of great experimental importance in electronic industry. The varied values of optical dielectric constants (ε_{∞}) and average energy gaps (Eg) of ionic and covalent crystals have helped in developing many components for photo elastic, photo conducting and solar cell devices. For the study of dielectric behaviour, a number of classical models have been prescribed. However, the average energy gap is a quantum concept, hence a good number of quantum theories are also prescribed. Most satisfactory results are obtained through the quantum ion dependent model^(1 to 5) established in last two decades. The model says that the dielectric behaviour of ionic crystals and partial ionic-covalent solids (with dominating ionic contribution) is completely cation dependent. The relation between ε_{∞} and the interionic separation R is suggested as⁽¹⁾

 $\varepsilon_{\infty} = 1 + BR^{S} \tag{1}$

where B is a characteristic constant of a particular cation and S is a family characteristic constant. This ion dependent model is well appreciated in the review article of $Moss^{(6)}$.

In the present paper we propose to mix the solids of rocksalt and zinc blende structures, specially in partial ionic- covalent family II-VI. Mixed solids in different proportions are taken up and their electronic dielectric constants are evaluated by using the ion dependent concept.

Theory and Method of Calculation

Any two crystals X and Y with same cation in binary II-VI family are mixed in different proportion λ_x and λ_y respectively. The interionic separation of the mixed crystal (R) is found through the relation

$$\mathbf{R}^3 = \lambda_x \, \mathbf{R}_x^3 + \lambda_y \, \mathbf{R}_y^3 \qquad (2)$$

where R_x and R_y are the interionic separations of the pure II-VI crystals X and Y respectively. This formula is in accordance with the concept of volume mixing^(7,8). Calculated values of R for mixed crystals in different proportions are reported in tables 1 to 6. Since we are mixing two same cation solids, values of characteristic constants B and S for same cations are directly taken up from the tables, earlier predicted for pure II-VI family through ion dependent formulation⁽¹⁾. Thus, by using the values of R obtained from eqn. (2), we predict the values of ε_{∞} for different mixed crystals through eqn. (1). Tables 1 to 6 show the calculated values of ε_{∞} through ion dependent calculations.

VOL- VIII ISSUE- XI NOVEMBER 2021 PEER REVIEW e-JOURNAL IMPACT FACTOR 7.149 ISSN 2349-638x

In order to check the validity of the results we call back the well established classical Clausius Mossotti relation for electronic dielectric constants

$$\frac{\varepsilon_{\infty} - 1}{\varepsilon_{\infty} + 2} = \frac{4\pi}{3} \frac{\alpha}{V}$$
(3)

where α and V are the polarisability and volume per ion pair of the crystal respectively.

Kamiyoshi and Niagra⁽⁹⁾ and Varotsos⁽¹⁰⁾ considered the polarisability of mixed crystal to be equal to the sum of polarisabilities of the pure crystals taken in the proportion of mixing ratio. Thus eqn. (3) in modified⁽¹¹⁾ as

$$\frac{\varepsilon_{\infty} - 1}{\varepsilon_{\infty} + 2} = \lambda_{x} \left[\frac{(\varepsilon_{\infty})_{x} - 1}{(\varepsilon_{\infty})_{x} + 2} \right] + \lambda_{y} \left[\frac{(\varepsilon_{\infty})_{y} - 1}{(\varepsilon_{\infty})_{y} + 2} \right]$$
(4)

using experimental values of $(\varepsilon_{\infty})_x$ and $(\varepsilon_{\infty})_y$ for pure II-VI crystals⁽¹⁾, ε_{∞} for mixed crystals are

evaluated from eqn (4). These values are also reported in table 1 to 6 for comparison.

 Table – 1 : Rocksalt Structure

MgS –	R	(€∞)a	(£∞)b	(ε∞) _{Exp.}
MgSe	R		(60)0	(Coo)Exp.
(% MgSe)				
0	2.60	5.09	5.09	5.1
10	2.61	5.16	5.16	-
20	2.63	5.22	5.25	-
30	2.64	5.29 🔪	5.32	SAT
40	2.65	5.35	5.38	-
50	2.67	5.42	5.46	-
60	2.68	5.48	5.56	-
70	2.69	5.54	5.63	W ail
80	2.70	5.61	5.71	-
90	2.72	5.67	5.79	-
100	2.73	5.74	5.89	5.9

Table – 2 : Rocksalt Structure

CaO – CaS (%CaS)	R	(€∞)a	(ɛ∞)b	(£∞)Exp.
0	2.40	3.21	3.30	3.3
10	2.45	3.36	3.40	-
20	2.50	3.50	3.50	-
30	2.55	3.65	3.61	-
40	2.59	3.79	3.71	-
50	2.64	3.94	3.84	-

60	2.68	4.08	3.96	-			
70	2.72	4.23	4.08	-			
80	2.76	4.37	4.21	-			
90	2.80	4.52	4.35	-			
100	2.84	4.66	4.49	4.5			
т	Table 3 · Dealscalt Structure						

Table – 3 : Rocksalt Structure

	SrSe –	R	$(\varepsilon_{\infty})_{a}$	(ε∞) _b	$(\varepsilon_{\infty})_{Exp.}$				
	SrTe								
	(% SrTe)								
	0	3.12	5.04	4.90	4.9				
	10	3.13	5.09	4.98	-				
1	20	3.14	5.14	5.06	-				
	5C30/j	3.16	5.18	5.14	-				
	40	3.17	5.23	5.23	-				
	50	3.18	5.28	5.32	-				
	60	3.19	5.33	5.41	-				
	70	3.20	5.38	5.50	-				
	80	3.22	5.42	5.59	-				
A REAL	90	3.23	5.47	5.69	-				
	100	3.24	5.52	5.79	5.8				
1									

Table – 4 : Zinc blende Structure

BeS –	R	(ε∞) _a	(ɛ∞)b	$(\varepsilon_{\infty})_{Exp.}$				
MeSe		-						
(% BeSe)		2						
0	2.10	6.48	7.09	7.1				
10	2.11	6.56	7.11	-				
20	2.12	6.65	7.13	-				
30	2.13	6.73	7.15	-				
40	2.14	6.81	7.17	-				
50 29	2.15	6.89	7.19	-				
60	2.16	6.97	7.21	-				
70	2.17	7.06	7.23	-				
80	2.18	7.14	7.25	-				
90	2.19	7.22	7.27	-				
100	2.20	7.30	7.29	7.3				

Table – 5 : Zinc blende Structure

ZnS –	R	$(\epsilon_{\infty})_a$	$(\epsilon_{\infty})_{b}$	(ε∞) _{Exp.}
ZnSe				
(% ZnSe)				
0	2.36	5.38	5.19	5.2
10	2.37	5.43	5.26	-
20	2.38	5.48	5.32	-
30	2.39	5.53	5.39	-
40	2.40	5.58	5.46	-
50	2.40	5.64	5.53	-
60	2.41	5.69	5.60	-

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	Aayushi International Interdisciplinary Research Journal (AIIRJ)								
١	/OL- VIII	ISSUE- XI	NOVE	MBER 2	2021	PEER REVIEW e-JOURNAL	IMPACT FACTOR 7.149	ISSN 2349-638x	
	70	2.42	5.74	5.67	-		$\varepsilon_{\infty} = n^2$ (5)		
	80	2.43	5.79	5.74	-	Thus	one can derive prope	er crystal mixture	
	90	2.44	5.85	5.82	-	for any def	finitely required refra	ctivity in opto-	
	100	2.45	5.90	5.89	5.9	electronic sys	stems.		

CdS –	R	(ε∞) _a	(ɛ∞)b	(ε∞) _{Exp.}
CdTe				
(% CdTe)				
0	2.52	5.53	5.19	5.2
10	2.55	5.68	5.35	-
20	2.58	5.84	5.52	-
30	2.60	5.99	5.70	
40	2.63	6.15	5.88	-
50	2.66	6.30	6.07	-
60	2.68	6.46	6.28	-
70	2.71	6.61	6.49	-
80	2.73	6.77	6.72	-
90	2.75	6.92	6.95	-
100	2.78	7.08	7.20	7.2

Analysis of Results

A comparison of $(\varepsilon_{\infty})_a$ evaluated from quantum ion dependent formulation and $(\varepsilon_{\infty})_b$ predicted from classical Clausius Mossotti relation in tables 1 to 6 shows excellent agreement. Close matching of the results from two entirely different approaches proves the reliability of the values of electronic dielectric constants of mixed crystals in II-VI binary family. Thus quantum ion dependent formulation can be used to predict the electronic dielectric constant of the mixed crystals in ionic I-VII solids as well as covalent III-V semiconductors also, since the basic cation anion dependent theory covers all subsets of binary A^N B^{8-N} family⁽¹⁾.

Other kinds of quantum approaches for the developments in the theory of mixed crystals have been considered, but only for pure ionic crystals⁽¹²⁾. Covalent semiconductors have not been touched at all till to date. The present approach through quantum ion dependent formulation can be well applied for covalent semiconductors also. In a wider arena, a huge range of values of ε_{∞} can be obtained by mixing the crystals in different proportions, which can be used for many different purposes. The electronic dielectric constant has a direct bearing on the optical refractive index (n) as

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